Structure of Bi₂Nd₄O₉ Monoclinic Phase

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A new phase, Bi₂Nd₄O₉, was synthesized in the binary system Bi₂O₃–Nd₂O₃ and structurally characterized from powder X-ray diffraction data. It crystallizes in monoclinic symmetry with lattice parameters a = 6.7037(2) Å, b = 3.9062(1) Å, c = 3.9542(1) Å, $\beta = 125.258(1)^{\circ}$, Z = 2, and a subcell space group of C2/m. Its average structure had been determined by X-ray powder diffraction and refined by the Rietveld method. The average structure can be described as an oxygen-deficient fluorite structure of the δ -Bi₂O₃ type with Nd and Bi randomly distributed on the cation site, similar to Bi₂La₄O₉ reported by Chen *et al.* (*J. Solid State Chem.* 124, 300 (1996)). In addition, a strong superstructure had been observed by transmission electron microscopy. © 2000 Academic Press

Key Words: Rietveld refinement; superstructure; Bi₂Nd₄O₉.

I. INTRODUCTION

The binary systems $Bi_2O_3-Ln_2O_3$ (Ln = Y, La-Lu) have been widely investigated over the Bi-rich region (1-7) because δ -Bi₂O₃ and doped Bi₂O₃ with fluorite-like structures exhibit high ionic conductivity. On the other hand, the *Ln*-rich side of the binary systems has also been explored (1,2,5-10) as the phases with high concentrations of *Ln* were expected to exhibit various types of magnetic ordering at low temperature (8). A large group of rhombohedral phases (Bi₃Ln₅O₁₂) (8, 9) and orthorhombic compounds (Bi₈La₁₀O₂₇) (11) has been characterized. Recently, Chen *et al.* (1) observed a new monoclinic phase, Bi₂La₄O₉, with the lattice parameters a = 6.8290(3) Å, b = 3.9887(1) Å, c = 4.0524(1) Å, $\beta = 125.094(3)^\circ$, and space group C2/m.

In the present work, a monoclinic compound $Bi_2Nd_4O_9$ was synthesized in the Nd-rich Bi_2O_3 -Nd₂O₃ system. Here we report X-ray diffraction data and the average structure of the monoclinic phase.

II. EXPERIMENTAL

The compound was prepared by solid-state reaction of mixtures of Nd_2O_3 (purity 99.99%) and Bi_2O_3 (purity

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99.99%). Before weighing, Nd_2O_3 was preheated at 900°C for 12 h to remove the carbon dioxide and water it absorbs in air. The mixture was ground, preheated at 800°C for 12 h in air, and then fired at 1100°C for 5–7 days with intermediate grindings. The sample was finally quenched at room temperature in air. The weight loss, presumably due to Bi_2O_3 vaporization during heating, was always less than 0.8% of the total sample weight (1).

Phase identification, determination of lattice parameters, and Rietveld refinement were performed on a Rigaku diffractometer using CuK α radiation at room temperature. For phase identification, the data were collected in the normal scanning mode from 10° to 80°. The lattice parameters were calculated with the program DICVOL91 (12). For Rietveld refinement, the data were collected in the step scanning mode with steps of 0.02° (2 θ). An angular range from $2\theta = 25^{\circ}$ to 130° and a measuring time of 1 s per step were applied. The divergence, scattering, and receiving slits were set at 1°, 1°, and 0.1 mm, respectively. A total of 5251 points with a maximum count of 17,193 and 83 independent reflections were measured.

The observed powder diffraction data of Bi₂Nd₄O₉ are given in Table 1. The main peaks can be indexed using a monoclinic subcell with lattice parameters a = 6.7037(2) Å, b = 3.9062(1) Å, c = 3.9542(1) Å, and $\beta = 125.258(1)^{\circ}$ (F(25) = 52.2). The observed condition for diffraction allows for the space group C2/m, C2, or Cm. The cell can be transformed to another setting with a = 5.474(1) Å, c = 3.954(1) Å, $\beta = 90.91(2)^{\circ}$ b = 3.906(5) Å, and (F(25) = 50.6), corresponding to the space groups I2/m, I2, or Im. It is evident that the cell parameters of its bodycentered cell are closely related to those of fluorite type δ -Bi₂O₃ (fcc) by the relations $b \approx c \approx a \sqrt{2}/2$. Since the unindexed diffraction peaks are weak, the approximate crystal structure could be refined with the subcell.

III. RIETVELD REFINEMENT

Because of the apparent close relationship with the fluorite structure, calculations were performed with starting



			Observed		Calculated					Observed		Calculated	
h	k	l	d	I/I_0	d	I/I_0	h	k	l	d	I/I_0	d	I/I_0
	*		4.605	2			4	0	$^{-2}$	1 501	11	1.591	0
	*		3.714	3			2	2	0	1.391	11	1.590	9
	*		3.573	5			2	2	-2	1.389	5	1.389	4
	*		3.412	1			4	0	0	1.368	2	1.369	2
0	0	1	3.232	31	3.229	31	2	0	-3	1.285	1	1.286	1
2	0	-1	3 184	100	3.182	100	1	1	2	1.277	4	1.277	
1	1	0	5.104	100	3.180	100	4	0	-3			1.277	4
	*		3.044	2			2	2	1	1.267	3	1.268	3
	*		2.904	2			1	3	0	1.266	3	1.267	3
1	1	-1	2.780	32	2.779	32	5	1	-2	1 250	(1.260	5
2	0	0	2.738	14	2.737	15	4	2	-1	1.259	0	1.259	
	*		2.011	2			3	1	-3	1.248	2	1.249	2
2	0	-2	1.977	9	1.977	7	0	2	2	1.244	2	1.244	1
1	1	1	1.961	19	1.961	19	1	3	-1	1.236	2	1.237	2
0	2	0	1.953	10	1.953	10	4	2	-2	1.233 4		1.234	4
3	1	-1	1.939	21	1.940	19	3	1	1		4	1.233	
1	1	-2	1.685	8	1.686	7	5	1	-1	1.222	2	1.222	2
3	1	-2	1 (71 1)	17	1.672	17	1	1	-3	1.142	2	1.142	1
0	2	1	1.0/1	17	1.671		2	0	2	1.133	1	1.133	1
2	0	1	1.666	4	1.667	4	5	1	-3	1.130	2	1.130	2
2	2	-1	1.664	9	1.665	9	1	3	1	1.129	2	1.129	2
3	1	0	1.653	8	1.653	9	3	3	-1	1.125	2	1.125	2
4	0	-1	1.647	4	1.647	4	4	2	0	1.120	2	1.121	2
0	0	2	1.614	3	1.614	3	6	0	-2	1.117	1	1.117	1

TABLE 1X-Ray Powder Diffraction Data of $Bi_2Nd_4O_9$ for Cu Radiation, $\lambda = 1.5406 \text{ Å}$

Note. For the indexed subcell, the Smith-Snyder figure of merit F(25) = 52.2 (0.0082, 58). Reflections marked by * are due to the superstructure.

positional parameters according to those of the classical fluorite. The subcell then should contain one $(Bi, Nd)_2O_3$. Starting with C2/m, the Bi and Nd atoms were located statistically on the site 2a(0,0,0) with the initial ratio Nd/Bi = 2.0. The oxygen atoms were set on the site 4i(x, 0, z) with x = 0.25 and z = 0.75. The occupancy factor for oxygen was set to 0.75, corresponding to three oxygen atoms in the cell.

Using the computer program DBWS-9411 (13), the Rietveld refinement was carried out. The pseudo-Voigt function was used for the simulation of the peak shapes (14). Intensities within 16 times of the full width at half-maximum were considered to contribute to the reflection. Peaks below 75° (2θ) were corrected for asymmetry effects after Rietveld (15). The preferred orientation was refined using the March-Dollase function. The background was refined by use of a third-order polynomial in 2θ . In the first step, the scattering factor, zero point, background, profile parameters, lattice constant, and atomic coordinates were fitted.

Under these conditions, successive refinements of the positional parameters of oxygen and the isotropic thermal parameters lead to $R_{wp} = 18.72\%$, $R_{exp} = 6.70\%$, and $R_{Bragg} = 7.73$. At this stage, refinement of the ratio Bi/Nd

while keeping their isotropic thermal parameters fixed confirmed the ratio Bi/Nd = 1:2.

Attempts with space groups C2 and Cm only yielded larger agreement factors. Hence, the most probable space group is C2/m. The results of the refinment are plotted in Figure 1. Table 2 gives the structure parameters of the final refinement. Figure 2 shows the average crystal structure of Bi₂Nd₄O₉.

IV. DISCUSSION

Figure 2 gives the monoclinic unit cell of $Bi_2Nd_4O_9$ that has been shown to be isostructural with $Bi_2La_4O_9$ (1). However, *a*-, *b*-, and *c*-axes shrank from 6.8290 Å, 3.9887 Å, and 4.0524 Å for $Bi_2La_4O_9$ (1) to 6.703 Å, 3.906 Å, and 3.954 Å for $Bi_2Nd_4O_9$. The unit cell volume of $Bi_2Nd_4O_9$ was found to be smaller than that of $Bi_2La_4O_9$. The bond length between Nd and O atoms is also shorter than that between La and O atoms in $Bi_2La_4O_9$. All these verities can be attributed to the smaller ionic radius of Nd^{3+} (1.08 Å) in comparison with that of La^{3+} (1.15 Å).

From Table 2, it is obvious that the oxygen atom site in $Bi_2Nd_4O_9$ tends to occupy the ideal site, (0.25, 0, 0.75), of the



represent the experimental values and solid lines the calculated pattern. The solid line at the bottom is the difference between the experimental and calculated values. The vertical markers show positions calculated for Bragg reflections.

fluorite structure compared with the case in $Bi_2La_4O_9$. At the same time, in the body-centered cell, $\beta = 91.30^{\circ}$ (1) in $Bi_2La_4O_9$ and $\beta = 90.91^\circ$ in $Bi_2Nd_4O_9$. The decrease of β indicates that the monoclinic structure is inclined to orthorhombic structure with the decrease of the ionic radius of Ln. Therefore, it is reasonable to presume that the smaller ionic radius of Ln is favorable for the oxygen site being at the ideal site and stable non-monoclinic compounds. Maybe this is why there are no other monoclinic isostructural compounds with further lanthanide ions. The monoclinic structure is no longer stable with the smaller cations.

The isothermal temperature factors B for (Bi,2Nd) and O are 1.8(3) Å² and 5.6(4) Å², respectively. The values are rather large. Comparably large B values were also observed $Bi_{0.7}La_{0.3}O_{1.5}$ [B(La,Bi) = 0.91–1.35 Å², B(O) = for 1.40–4.93 Å²] (16), $Bi_8La_{10}O_{27}$ [*B*(La,Bi) = 1.6–2.5 Å², $B(O) = 1.9-6.3 \text{ Å}^2$ (11), and $Bi_2La_4O_9$ [$B(La,Bi) = 1.8 \text{ Å}^2$, $B(O) = 1.5 \text{ Å}^2$ (1). The large B values are usually attributed to the ordering phenomena which will result in the superstructure in the compounds. Because of the better ordering of atoms in $Bi_2Nd_4O_9$, B values are larger than those in

TABLE 2 Final Rietveld Refinement Parameters of Bi₂Nd₄O₉ from X-Ray **Powder Diffraction Data**

Atom	Site	x/a	y/b	z/c	$B~(\mathrm{\AA}^2)$	Occupancy
Nd) Bi	2 <i>a</i>	0	0	0	1.8(1)	0.667
0	4 <i>i</i>	0.264(6)	0	0.774(5)	5.6(4)	0.333 0.75

Note. $R_{wp} = 18.72\%$, $R_{exp} = 6.70\%$, $R_{Bragg} = 7.73$.



FIG. 2. Monoclinic subcell of Bi₂Nd₄O₉. The open circles stand for Nd or Bi atoms and the filled circles for oxygen.

other compounds. Correspondingly, transmission electron microscopy investigation revealed the existence of a strong superstructure in $Bi_2Nd_4O_9$.

Figure 3 shows the electron diffraction pattern along [100] of the body-centered cell. Strong spots originate from



FIG. 3. Electron diffraction pattern of Bi₂Nd₄O₉. The electron beam is parallel to [100]. The superstructure is along [001] with a period 3c (horizontal direction).

8000

the average structure, while weak spots indicate a possible triple superstructure along [001] of the body-centered cell. The weak spots are more noticeable comparing the reported data. Thus TEM indicates a strong superstructure in the compound.

A weak superstructure has also been observed in $Bi_3Nd_5O_{12}$ (8), $BiLa_2O_{4.5}$ (9), $Bi_2La_4O_9$ (1), $Bi_3Pr_7O_x$ (6), $Bi_2Gd_3O_{7.5}$ (5), etc. The superstructure generally originates from atomic ordering in compounds. In the above-mentioned compounds the superstructural spots were usually very weak and *B* values were smaller than 2.0. However, in the $Bi_2Nd_4O_9$ compound, the superstructural spots were very noticeable. The strong superstructure should originate from the better order of the Bi/Nd atoms in $Bi_2Nd_4O_9$ and results in the larger B(Bi,Nd) and B(O). At the same time, the superstructural lines which cannot be indexed by C2/m resulted in the larger R_{wp} refined by the Rietveld method.

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REFERENCES

- 1. X. L. Chen, W. Eysel, and J. Q. Li, J. Solid State Chem. 124, 300 (1996).
- 2. E. M. Levin and R. S. Roth, J. Res. Nat. Bur. Stand. 68, 197 (1964).
- 3. H. Iwahara, T. Esaka, and T. Sato, J. Solid State Chem. 39, 173 (1981).
- 4. A. Watababe, Solid State Ionics 34, 35 (1989).
- J. L. Yang, J. K-. Liang, W. H. Tang, Y. Shi, X. L. Chen, G. H. Rao, J. Solid State Chem. 125, 85 (1996).
- J. L. Yang, J. K. Liang, W. H. Tang, Y. Shi, X. L. Chen, and G. H. Rao, J. Alloys Compds. 252, 143 (1997).
- X. L. Chen, F. F. Zhang, Y. M. Shen, J. K. Liang, W. H. Tang, and Q. Y. Tu, J. Solid State Chem. 139, 398 (1998).
- R. Horyn, M. Wolcyrz, and A. Wojakowski, J. Solid State Chem. 116, 68 (1995).
- 9. M. Wolcyrz, L. Kepinski, and R. Horyn, J. Solid State Chem. 116, 72 (1995).
- Y. Z. Sun, G. H. Rao, J. L. Liang, W. H. Tang, X. L. Chen, and J. K. Liang, J. Alloys Compds. 248, 106 (1997).
- C. Michel, V. Caignaert, and B. Raveau, J. Solid State Chem. 90, 296 (1991).
- 12. A. Boultif and D. Louer, J. Appl. Crystallogr. 24, 987 (1991).
- R. A. Young, A. Sakthivel, T. S. Moss, and C. O. Paiva-Santos, J. Appl. Crystallogr. 28, 366(1995).
- 14. R. A. Young and D. B. Wiles, J. Appl. Crystallogr. 15, 430 (1982).
- 15. H. M. Rietveld, J. Appl. Crystallogr. 2, 65 (1969).
- D. Mercurio, M. Ei Farissi, J. C. Champarnaud-Mesjard, and B. Frit, J. Solid State Chem. 80, 133 (1989).